

# A SYMPHONY OF GIANT MONSTERS: SEARCHING FOR INTERMEDIATE-MASS BLACK HOLE BINARIES IN ADVANCED LIGO AND ADVANCED VIRGO'S THIRD OBSERVING RUN

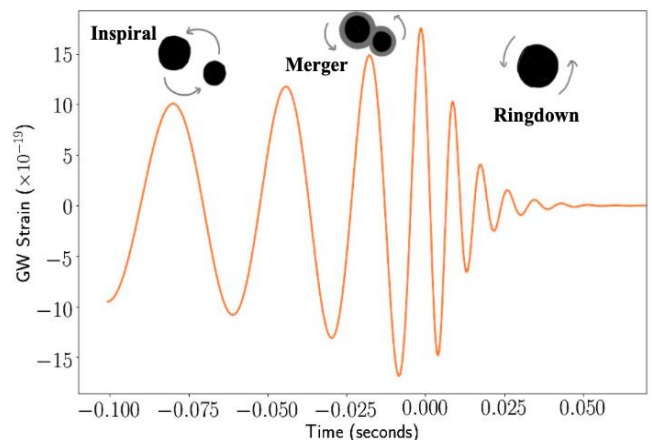
The search for intermediate-mass black holes ([IMBHs](#)) has been a focus of astronomers' attention for many years. The Advanced [LIGO](#) and Advanced [Virgo](#) detectors completed their first two observing runs (O1 and O2) between 2015 and 2017, and [no significant detections of IMBH mergers were made](#). The first half of our [third observing run](#) (O3) marked the first confirmed detection of an IMBH remnant in a binary system: [GW190521](#), a discovery which we announced in 2020. In pursuit of finding more such sources, we present here a dedicated IMBH search of the data from O3. In this new search we again found [GW190521](#), but with **higher significance**. We also found some other possible IMBH candidates, but none qualify as significant. However, the results of our search have improved our understanding of the astrophysical merger rates of IMBH binaries.

## IMBHs: THE MYSTERIOUS BRIDGE BETWEEN STELLAR-MASS AND SUPERMASSIVE BLACK HOLES

In our Universe, black holes (BHs) are broadly divided into three categories based on their masses. These are stellar-mass BHs, supermassive black holes (SMBHs), and IMBHs. The lightest among the three are stellar-mass BHs. These have masses up to 100 times that of the Sun ( $M_{\odot}$ ) and are formed by the explosion of a dying star of mass above  $20 M_{\odot}$ . To this date, the advanced LIGO-Virgo detectors have observed  $\sim 50$  gravitational waves (GW) signals from coalescing compact binaries, published in the catalogs [GWTC-1](#) and [GWTC-2](#).

SMBHs, on the other hand, have masses above  $\sim 100,000 M_{\odot}$ . Several theories suggest that these BHs form as a result of multiple collisions of massive BHs or IMBHs. In dense environments, the end products of first-generation mergers of compact objects can undergo further collisions to form massive second-generation BHs. This process could extend to further generations, eventually leading to what is referred to as a **hierarchical merger scenario**. So far, [observations](#) in the electromagnetic (EM) spectrum have shown that SMBHs reside at the center of most galaxies in the Universe. Although there have been shreds of evidence of SMBH mergers from EM observations, no observations are expected in the LIGO-Virgo frequency band – although the future space-based mission, [LISA](#), has excellent prospects to detect these SMBH mergers to very great cosmic distances.

As their name suggests, the third category, IMBHs, lies in between stellar-mass BHs and SMBHs. Their masses range from  $\sim 100$  to  $100,000 M_{\odot}$ . Some theories suggest that these BHs are the remnants of the direct core-collapse of massive first-generation, [low-metallicity Population III](#) stars, and others suggest that they are formed in the hierarchical mergers of stars in dense young star clusters. Either way, their very existence remains a puzzle until one makes a direct observation. At present, [HLX-1](#) is the strongest IMBH candidate seen in the X-ray part of the EM spectrum, with an inferred mass of  $\sim 3000$  to  $300,000 M_{\odot}$ . IMBHs can be directly observed in the GW window with Advanced LIGO and Virgo, however, by detecting the GWs emitted from a binary merger involving an IMBH – either as a binary component or as the merger remnant.



*Figure 1: A pictorial representation of the three phases of a GW signal from a massive BH binary. The quantity on the vertical axis represents the dimensionless GW strain (deformation of space-time) and the horizontal axis shows the duration of the signal.*

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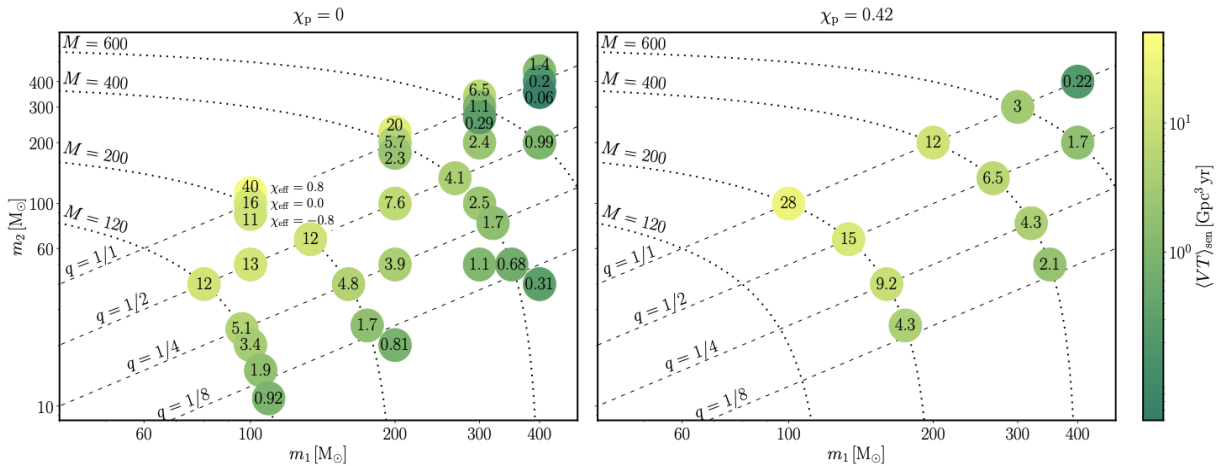
## WHY IS AN IMBH SEARCH IMPORTANT?

IMBH systems serve as a perfect laboratory to test the [theory of general relativity \(GR\)](#). Some of these binaries could have highly unequal, or *asymmetric*, masses. In general, such asymmetric binaries can produce GW [higher harmonics](#) along with their fundamental tone, just like what we see in musical instruments. These higher harmonics are also referred to as higher-order modes. One of our recent LIGO-Virgo detections, [GW190412](#), showed the [existence of these higher-order modes](#): this was the first time that scientists had detected a perfect [fifth](#) in the symphony of merging BHs. Detecting such asymmetric binaries can widen our understanding of the fundamental properties of BHs such as their mass, spin, and [charge](#) – as described by [the no-hair theorem](#). Another advantage of studying IMBH binaries is that the results can inform us about the stellar evolution of massive stars. Can these studies also provide insight into the evolutionary link between stellar-mass BHs and SMBHs? Well, the LIGO and Virgo detectors are insensitive to GWs from the even heavier SMBHs, so to obtain a more complete picture we may have to wait for LISA observations of SMBH mergers.

## HOW DO WE SEARCH FOR IMBH BINARIES?

GWs produced from the coalescences of compact binaries like neutron stars or BHs stretch and squeeze the [space-time](#) around interferometric detectors. The deformations created by this stretching and squeezing are precisely measured by the [LIGO](#) and [Virgo](#) detectors via [laser interferometry](#). The detectors not only capture GW signals but also see many types of noise coming from terrestrial and instrumental sources – including some which can be much “louder” than the GW signals themselves. To find astrophysical GW signals embedded in these noisy data, scientists use special techniques such as “[matched filtering](#)”, which uses the GW signal model predicted from GR to look for similar patterns in the data. Another [model-independent](#) technique searches for excess power across the detectors occurring within the light travel time between them – meaning that this excess could plausibly be coming from a GW signal.

A generic GW signal model is represented by its [inspiral](#), [merger](#), and [ringdown](#) phases (see [Figure 1](#)), the durations of which depend on the properties of the GW source. The heavier the binary, the shorter the GW signal emitted from it. Extracting such short signals from noisy data is indeed challenging, and the difficulty level can increase further due to terrestrial or instrumental artifacts camouflaging the actual signal. Therefore, it becomes imperative to carefully tune our search techniques so that they can best identify only real GW signals.



**Figure 2:** (Figure 1 from the paper) The averaged sensitive volume-time  $\langle VT \rangle_{\text{sen}}$ , in  $\text{Gpc}^3 \text{yr}$  (in gigaparsecs cubed years) for the targeted non-precessing (left panel) and precessing (right panel) IMBH binaries. Each circle corresponds to one class of IMBH binary with the given masses. For the non-precessing class of IMBH binaries, the sensitive reach of our IMBH search decreases with increasing total mass ( $M$ ) of the binary system for a constant mass ratio ( $q$ ). In the case of precessing binaries, indicated by non-zero values of  $\chi_p$ , the search sensitivity is directly correlated with how positive is the total effective spin of the system.

In the O3 observing run, a combination of model-based matched filtering and an unmodeled search was used to look for IMBH binaries. The combination of the two methods results in a list of plausible IMBH binary candidates that are found in more than one detector. Among these candidates, we select the ones that have a remnant mass above  $100 M_{\odot}$  and the primary mass (i.e. the mass of the heavier BH in the binary) above  $65 M_{\odot}$ , which is the lower bound of the [pair-instability mass gap](#). We also require that our selected candidates have low values of [false alarm rate](#) (FAR), which means that the [odds](#) are small of finding a candidate with similar loudness that is simply background noise.

## WHAT DID WE FIND?

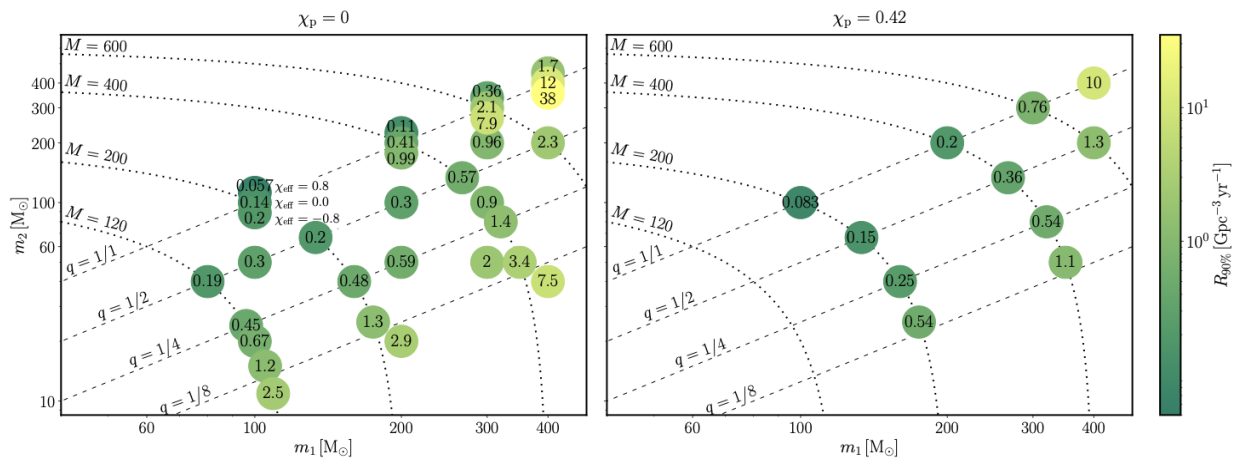
Based on our search selection criteria, we find **GW190521** to be the **most significant IMBH** candidate in the combined search. The discovery of GW190521 was announced on September 2, 2020, and the signal properties were studied in a [dedicated publication](#). The event originated from a merger of two BHs weighing  $66 M_{\odot}$  and  $85 M_{\odot}$ , leaving behind a remnant BH of mass  $\sim 142 M_{\odot}$ . In this search, we detect this event with higher significance than the previous analysis.

The IMBH search being sensitive to high-mass binaries makes it optimal for recovering heavy binary black hole (BBH) systems. With this search, we also recover GW signals from other heavy BBH sources seen in GWTC-2. However, there were no new significant IMBH signals observed in this search.

## WHAT CAN WE SAY ABOUT THE MERGER RATE OF IMBH BINARIES?

Estimating how often IMBH binaries merge in a given volume of the Universe and a given time (their so-called “merger rate density”) requires us to understand how sensitive our search is – i.e. how likely we are to detect signals from IMBH binaries. To investigate this, we perform simulations that involve “injecting” artificial GW signals into the data. First we use [Numerical Relativity](#) models that generate GW signals by numerically solving [Einstein equations](#) for coalescing IMBH binaries with total masses up to  $800 M_{\odot}$  and mass ratios (denoted by the parameter  $q$ ) varying between 1-to-1 and 1-to-10. We then distribute these signals uniformly in terms of their sky location, binary [orientation](#), and [comoving volume](#) – and use our combined search to try to recover them from the data. The recovered signals in a given volume and observational time provides us with a measure of the sensitivity of our search (see [Figure 2](#)) and places bounds on the merger rates of IMBH binary sources, as shown in [Figure 3](#).

From our studies, we found that the merger rate for equal-mass IMBH binaries with a total mass of  $200 M_{\odot}$  and spins of 0.8 is 0.06 per gigaparsec cubed per year. This upper limit is 3.5 times lower than the value reported in [the previous IMBH search from the O1 and O2 observing runs](#). This rate translates into 0.2 mergers over the typical lifetime of a [globular cluster](#) (assumed to be 10 billion years). We also revise the merger rate for GW190521-like sources to 0.08 per gigaparsec cubed per year.



**Figure 3:** (Figure 3 from the paper) The 90% upper limit on the merger rate density ( $R_{90\%}$ ) in  $\text{Gpc}^{-3} \text{yr}^{-1}$  (per gigaparsec cubed and per year) for the targeted non-precessing (left panel) and precessing (right panel) IMBH binaries. Again each circle corresponds to one class of IMBH binary with the given masses.

## GLOSSARY

**Higher-order-modes:** Sub-dominant terms in the [spherical harmonics](#) expansion that describes GWs beyond the dominant, quadrupolar term. The quadrupole term is analogous to the fundamental tone in sound waves.

**Mass ratio:** Ratio of the lighter component’s mass to the heavier component’s mass ( $q = m_2/m_1$ ).

**Gigaparsec (Gpc):** A measure of distance equal to one billion [parsec](#), which is about three billion light years or about  $3 \times 10^{22}$  km.

**Precessing binaries:** Binary systems where the components’ spin direction is misaligned with the direction of their orbital [angular momentum](#).  $\chi_{\text{eff}}$ , called the “effective aligned spin” is a parameter that characterizes a particular combination of the compact objects’ spins before they merge. A negative value of  $\chi_{\text{eff}}$  indicates anti-aligned spins with respect to the angular momentum of the system.  $\chi_p$ , called the “effective precession spin”, is another combination of individual spins that measures how much [precession](#) is present in the system, with a higher value indicating more precession.

**Sensitive volume-time:** The space-time volume within which the search technique is sensitive to detecting GW sources.

**Solar mass:** Mass of the Sun ( $M_{\odot}$ ), equal to  $1.989 \times 10^{30}$  kg.

**Upper limit:** A statement about the maximum value a physical quantity is likely to have given the results of our search.

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